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# Statistical Properties of Wave Chaotic Scattering and Impedance Matrices

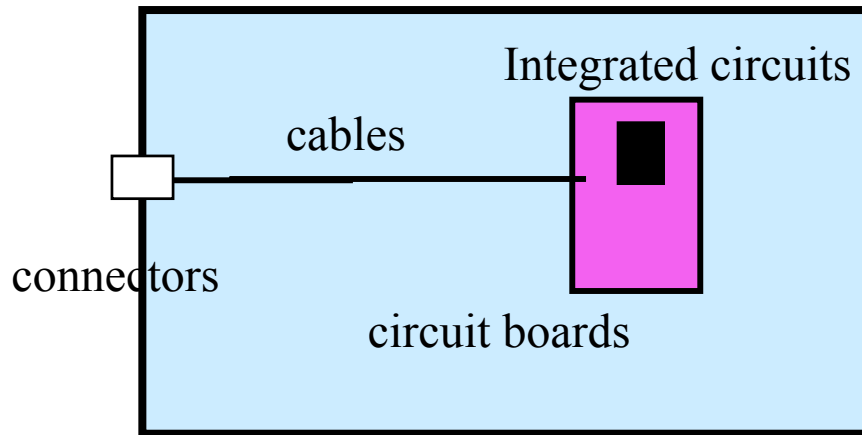
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MURI Students: Xing Zheng, Sameer Hemmady  
Collaborators: Shmuel Fishman, Richard Prange

AFOSR-MURI Program Review



# Electromagnetic Coupling in Computer Circuits

## Schematic



- What can be said about coupling without solving in detail the complicated EM problem ?
- **Statistical Description !**  
(Statistical Electromagnetics, Holland and St. John)

- Coupling of external radiation to computer circuits is a complex processes:

apertures  
resonant cavities  
transmission lines  
circuit elements

- Intermediate frequency range involves many interacting resonances
- System size  $\gg$   
Wavelength
- Chaotic Ray Trajectories

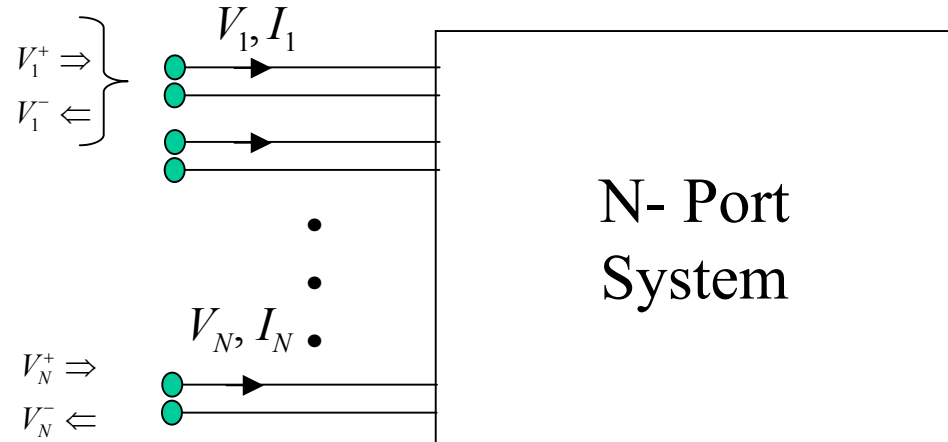


# Z and S-Matrices

## What is $S_{ij}$ ?

N ports

- voltages and currents,
- incoming and outgoing waves



Z matrix

$$\begin{pmatrix} V_1 \\ V_2 \\ \vdots \\ V_N \end{pmatrix} = \mathbf{Z} \begin{pmatrix} I_1 \\ I_2 \\ \vdots \\ I_N \end{pmatrix}$$

voltage

current

S matrix

$$\begin{pmatrix} V_1^- \\ V_2^- \\ \vdots \\ V_{N1}^- \end{pmatrix} = \mathbf{S} \begin{pmatrix} V_1^+ \\ V_2^+ \\ \vdots \\ V_{N1}^+ \end{pmatrix}$$

outgoing

incoming

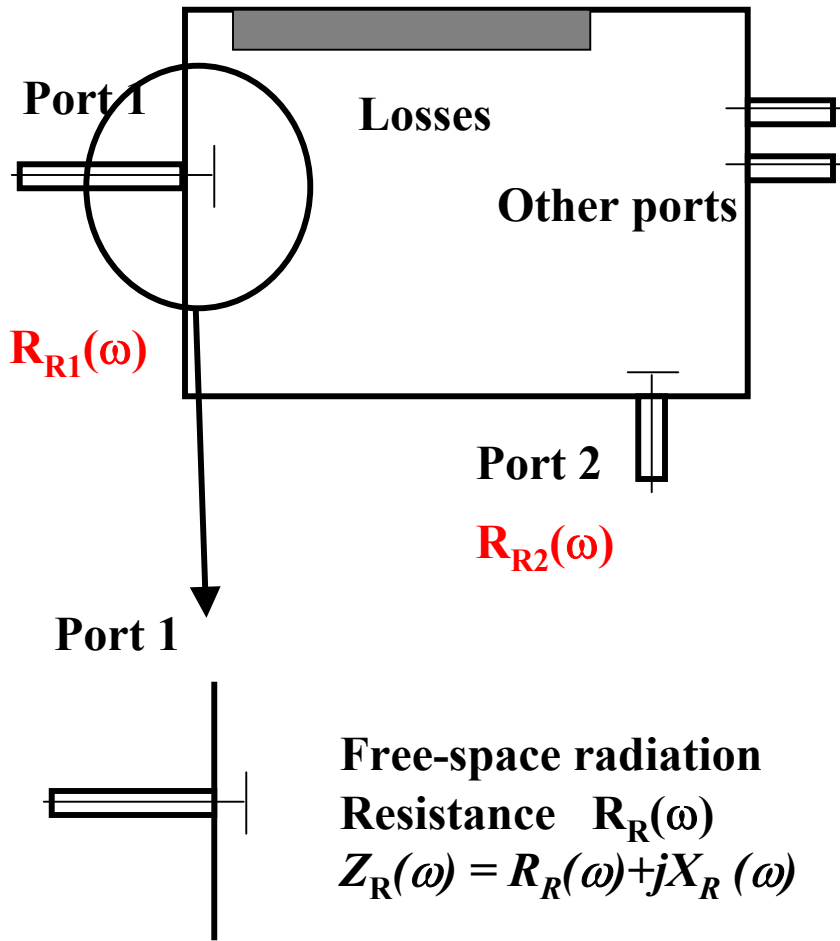
$$\mathbf{S} = (\mathbf{Z} + \mathbf{Z}_0)^{-1} (\mathbf{Z} - \mathbf{Z}_0)$$

$$\mathbf{Z}(\omega), \mathbf{S}(\omega)$$

- Complicated function of frequency
- Details depend sensitively on unknown parameters



# Statistical Model of Z Matrix



## Statistical Model Impedance

$$Z_{ij}(\omega) = -\frac{j}{\pi} \sum_n R_{Ri}^{1/2}(\omega_n) R_{Rj}^{1/2}(\omega_n) \frac{\Delta\omega_n^2 w_{in} w_{jn}}{\omega^2 (1 + jQ^{-1}) - \omega_n^2}$$

System parameters

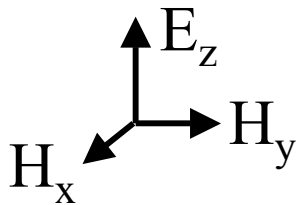
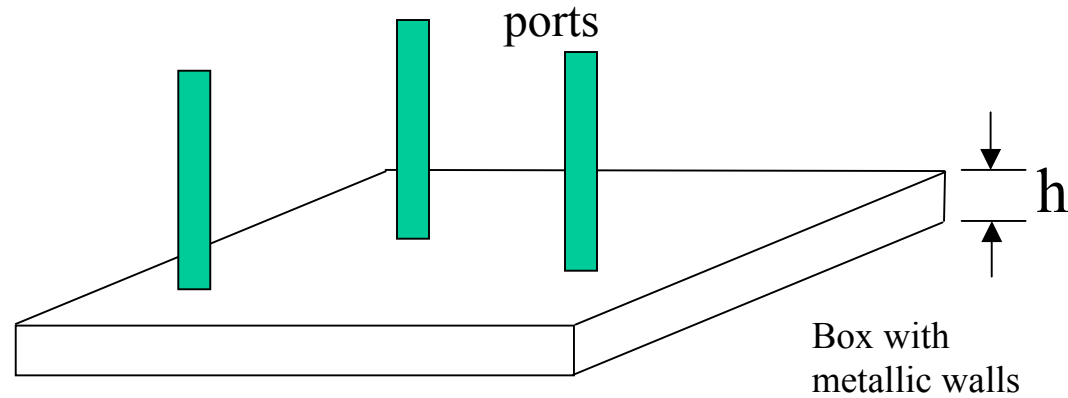
- Radiation Resistance  $R_{Ri}(\omega)$
- $\Delta\omega_n^2$  - mean spectral spacing
- $Q$  - quality factor

Statistical parameters

- $\omega_n$  - random spectrum
- $w_{in}$  - Gaussian Random variables



# Two Dimensional Resonators



Only transverse magnetic (TM) propagate for  $f < c/2h$

- Anlage Experiments
- Power plane of microcircuit

Voltage on top plate

$$E_z(x, y) = -V_T(x, y) / h$$



# Wave Equation for 2D Cavity $E_z = -V_T/h$

- Cavity fields driven by currents at ports, (assume  $e^{j\omega t}$  dependence) :

$$\nabla_{\perp}^2 V_T + k^2 V_T = -jkh \eta \sum_{\text{ports}} u_i I_i \quad \begin{array}{l} \eta = \sqrt{\mu / \epsilon} \\ k = \omega / c \end{array}$$

- Voltage at  $j^{\text{th}}$  port:  $V_j = \int dx dy u_j V_T$  Profile of excitation current

- Impedance matrix  $Z_{ij}(k)$ :  $Z_{ij} = -jkh \eta \int u_i \left( \nabla_{\perp}^2 + k^2 \right)^{-1} u_j dx dy$

- Scattering matrix:  $S(k) = (\mathbf{Z} + Z_0 \mathbf{I})^{-1} (\mathbf{Z} - Z_0 \mathbf{I})$



## Five Different Methods of Solution

Problem, find:  $Z_{ij} = -jkh\eta \int u_i (\nabla_{\perp}^2 + k^2)^{-1} u_j dx dy$

1. **Computational EM** - HFSS
2. **Experiment** - Anlage, Hemmady
3. **Random Matrix Theory** - replace wave equation with a matrix with random elements  
No losses
4. **Random Coupling Model** - expand in Chaotic Eigenfunctions
5. **Geometric Optics** - Superposition of contributions from different ray paths  
Not done yet

**Preprint available: Z and S 1.pdf**



## Expand $V_T$ in Eigenfunctions of Closed Cavity

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$$Z_{ij} = -jkh\eta \int u_i \left( \nabla_{\perp}^2 + k^2 \right)^{-1} u_j \, dx dy = -jkh\eta \sum_{\text{modes}-n} \frac{\langle u_i \phi_n \rangle \langle \phi_n u_j \rangle}{k^2 - k_n^2}$$

Where:

1.  $\phi_n$  are eigenfunctions of closed cavity  $\langle u_j \phi_n \rangle = \int dx dy u_j \phi_n$

2.  $k_n^2$  are corresponding eigenvalues

$$k_n = \omega_n / c$$

**$Z_{ij}$  - Formally exact**





# Random Coupling Model

## Replace $\phi_n$ by Chaotic Eigenfunctions

$$Z_{ij} = -jkh\eta \sum_{\text{modes}-n} \frac{\langle u_i \phi_n \rangle \langle \phi_n u_j \rangle}{k^2 - k_n^2}$$

1. Replace eigenfunction with superposition of random plane waves

$$\phi_n = \lim_{N \rightarrow \infty} \text{Re} \left\{ \sqrt{\frac{2}{AN}} \sum_{k=1}^N a_k \exp[i(k_n \mathbf{e}_k \cdot \mathbf{x} + \theta_k)] \right\}$$

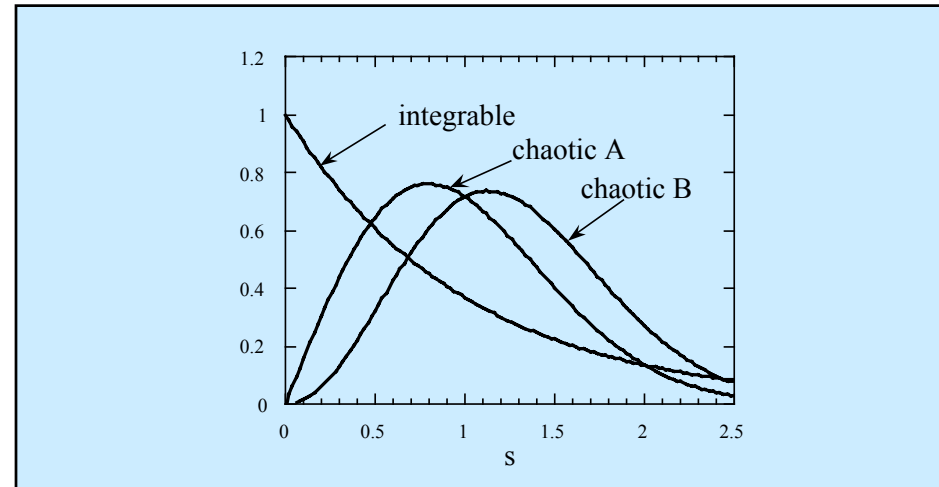
Random amplitude

Random direction

Random phase

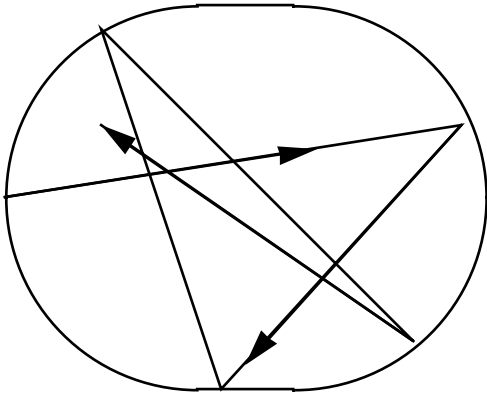
2. Eigenvalues  $k_n^2$  are distributed according to appropriate statistics

Normalized Spacing  $s_n = (k_{n+1}^2 - k_n^2) / \langle \Delta k^2 \rangle$





# Chaotic Eigenfunctions



Rays ergodically fill phase space.

Eigenfunctions appear to be a superposition of plane wave with random amplitudes and phases.

$$\phi_n = \lim_{N \rightarrow \infty} \operatorname{Re} \left\{ \sqrt{\frac{2}{AN}} \sum_{j=1}^N a_j \exp[i(k_j \cdot x + \theta_j)] \right\} \quad \text{Time reversal symmetry}$$

$$\phi_n = \lim_{N \rightarrow \infty} \frac{1}{\sqrt{2AN}} \sum_{j=1}^N a_j \exp[i(k_j \cdot x + \alpha_j)] \quad \text{Time reversal symmetry broken}$$

$k_j$  uniformly distributed on a circle  $|k_j| = k_n$

$\phi_n$  is a Gaussian random variable

TRS

$$P(\phi) \approx \exp[-|\phi|^2 / 2 \langle |\phi|^2 \rangle]$$



# Statistical Model for Impedance Matrix

$$Z_{ij}(k) = -\frac{j}{\pi} \sum_n R_{Ri}^{1/2}(k_n) R_{Rj}^{1/2}(k_n) \frac{\Delta k_n^2 w_{in} w_{jn}}{k^2(1 - jQ^{-1}) - k_n^2}$$

## System parameters

$$R_{Ri}(k) = \frac{k\eta}{4} \oint \frac{d\theta_k}{2\pi} |\bar{u}_i(k)|^2 = \text{Re}\{Z_{Ri}\}$$

-Radiation resistance for port i

$\Delta k_n^2 = 1/(4A)$  - mean spectral spacing

$Q$  -quality factor

## Statistical parameters

$w_{in}$  - Gaussian Random variables

$k_n$  - random spectrum



## Predicted Properties of $Z_{ij}$

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- Mean and fluctuating parts:  $Z_{ij} = \langle Z_{ij} \rangle + \delta Z_{ij}$

- Mean part:  $\langle Z_{ii} \rangle = jX_{R,i}(k)$  Radiation reactance  
(no losses)  $X_{R,i}(k) = \text{Im} \{ Z_R \}$

$$\langle Z_{ij} \rangle = 0, \quad i \neq j$$

- Fluctuating part: Lorentzian distribution  
-width radiation resistance  $R_{Ri}$

$$P(X_{ii}) = \frac{R_{Ri}}{\pi(X_{ii}^2 + R_{Ri}^2)}$$



# Numerical Test of $Z_{ii}$

**Numerically Generated Reactance**  $\xi = -\frac{1}{\pi} \sum_{n=1}^N \frac{w_n^2}{k^2 - k_n^2}$   $\xi = \bar{\xi} + \delta\xi$

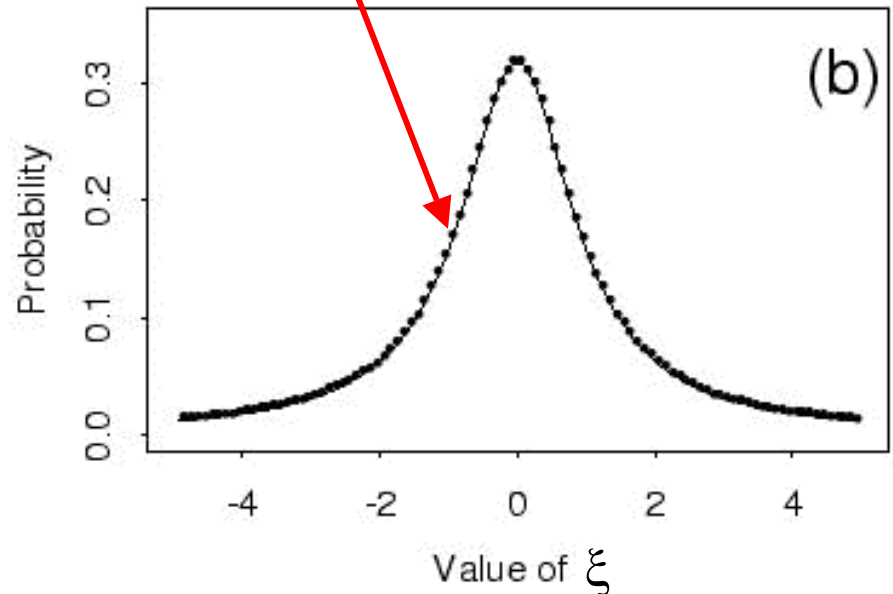
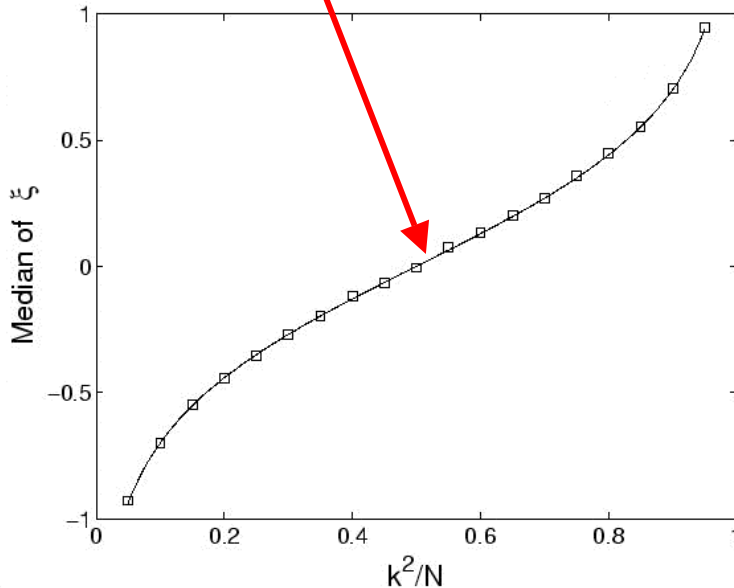
**Mean**  $\bar{\xi}$

**Theory:**  $\bar{\xi} = \frac{1}{\pi} \ln\left(\frac{N - k^2}{k^2}\right)$

**Fluctuation**  $\delta\xi$

$$P(\delta\xi) = \frac{1}{\pi(1 + \delta\xi^2)}$$

**10<sup>6</sup> realizations, N=2000,**

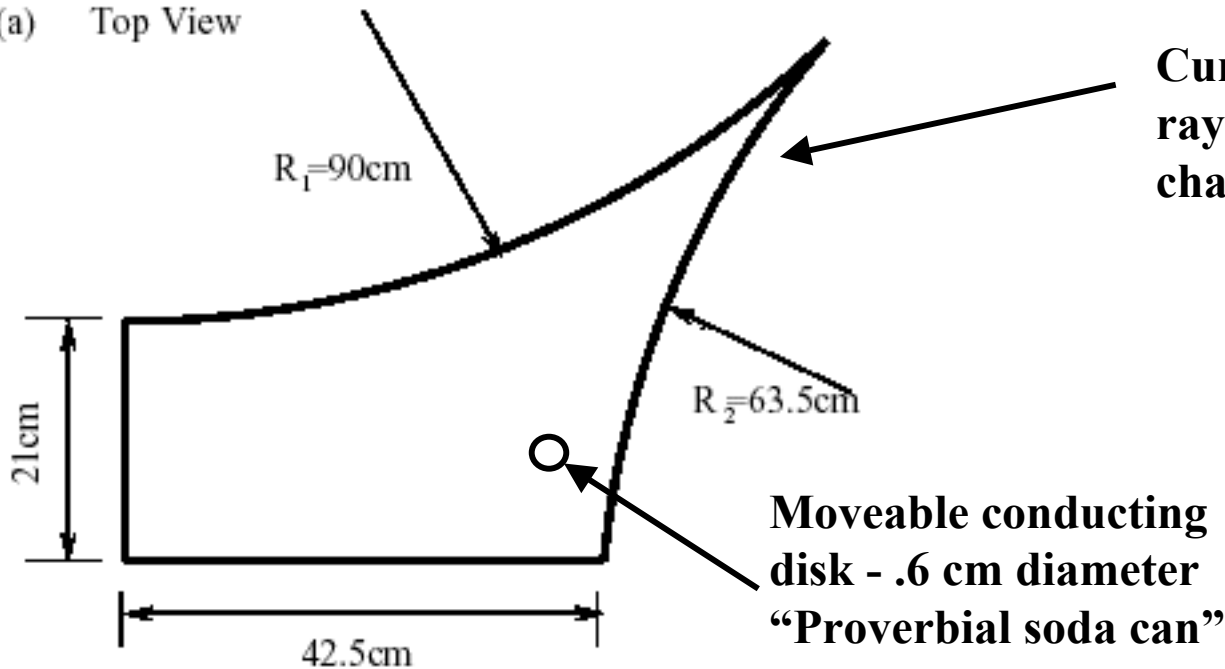




# HFSS - Solutions

## Bow-Tie Cavity

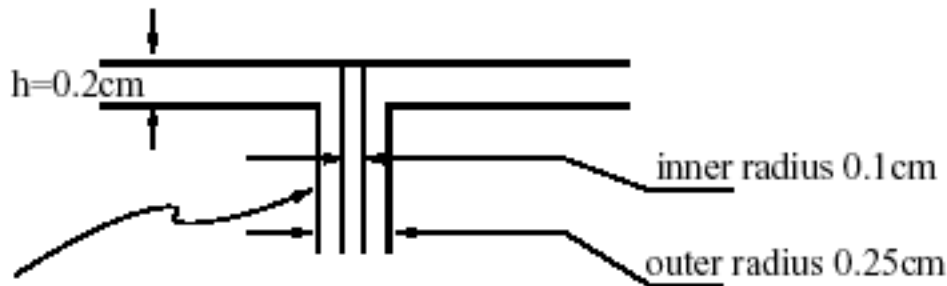
(a) Top View



Curved walls guarantee all ray trajectories are chaotic

Moveable conducting disk - .6 cm diameter  
"Proverbial soda can"

(b) Side View



Cavity impedance calculated for  
100 locations of disk  
4000 frequencies  
6.75 GHz to 8.75 GHz

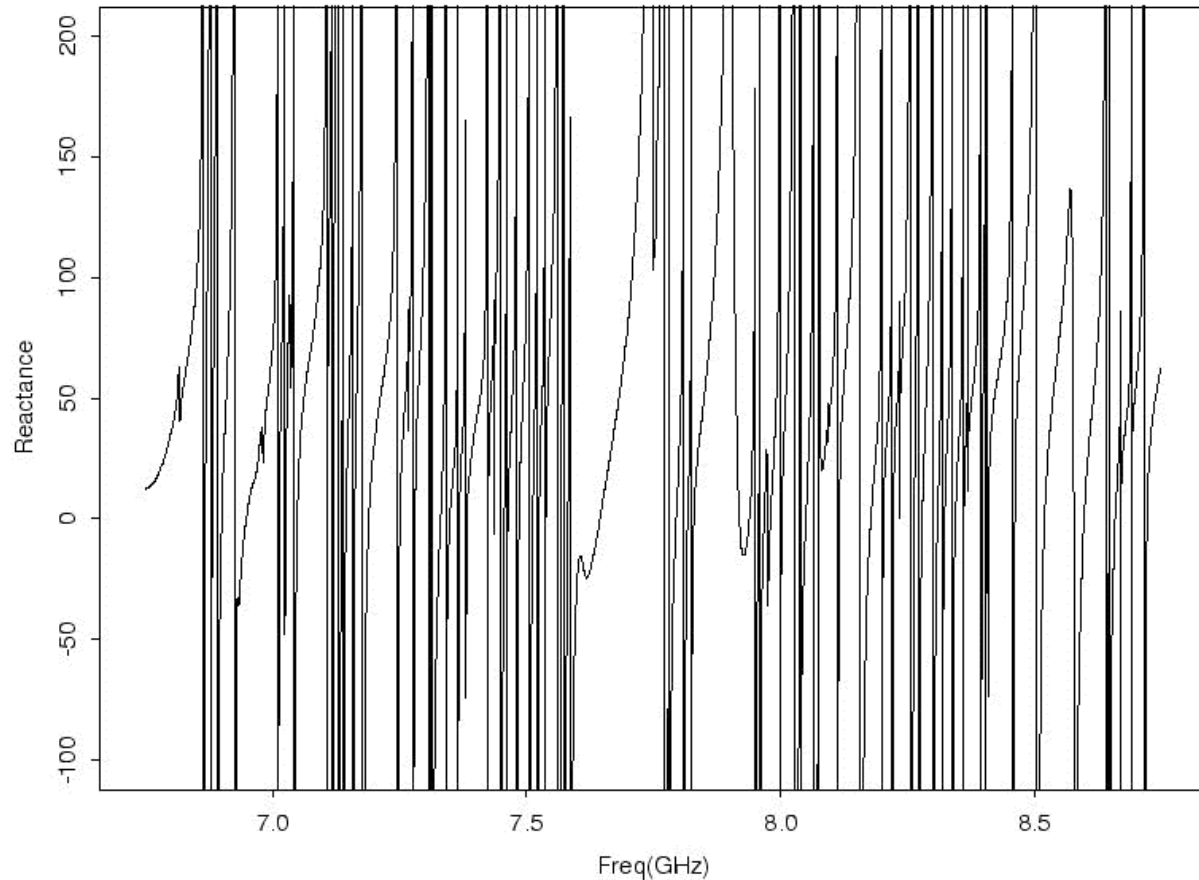


# Frequency Dependence of Reactance for a Single Realization

Mean spacing  $\delta f \approx .016$  GHz

$\Omega$

$$Z_{\text{cav}} = jX_{\text{cav}}$$





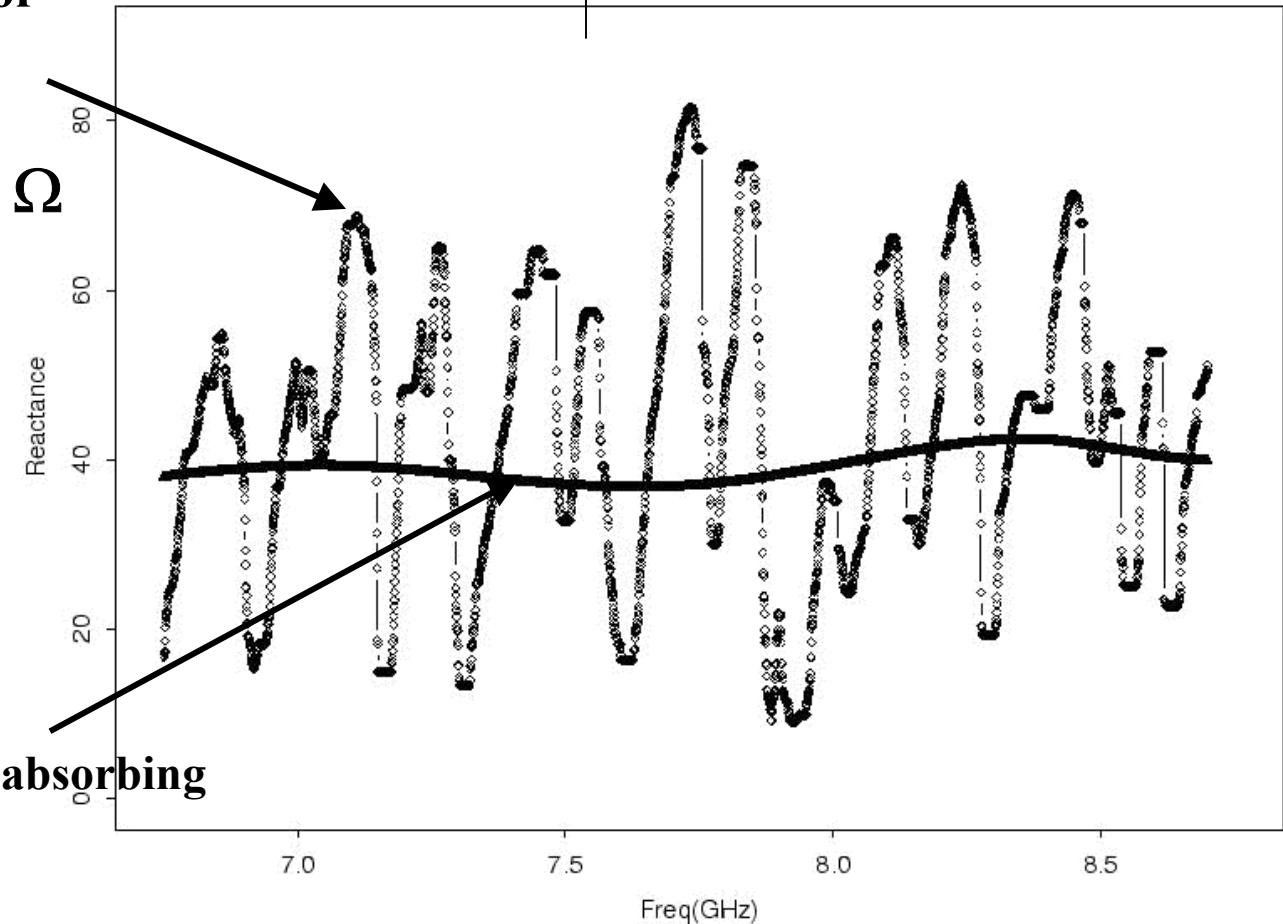
# Frequency Dependence of Median Cavity Reactance

$\Delta f = .3 \text{ GHz}, L = 100 \text{ cm}$

Median Impedance for 100 locations of disc

Effect of strong Reflections ?

Radiation Reactance HFSS with perfectly absorbing Boundary conditions



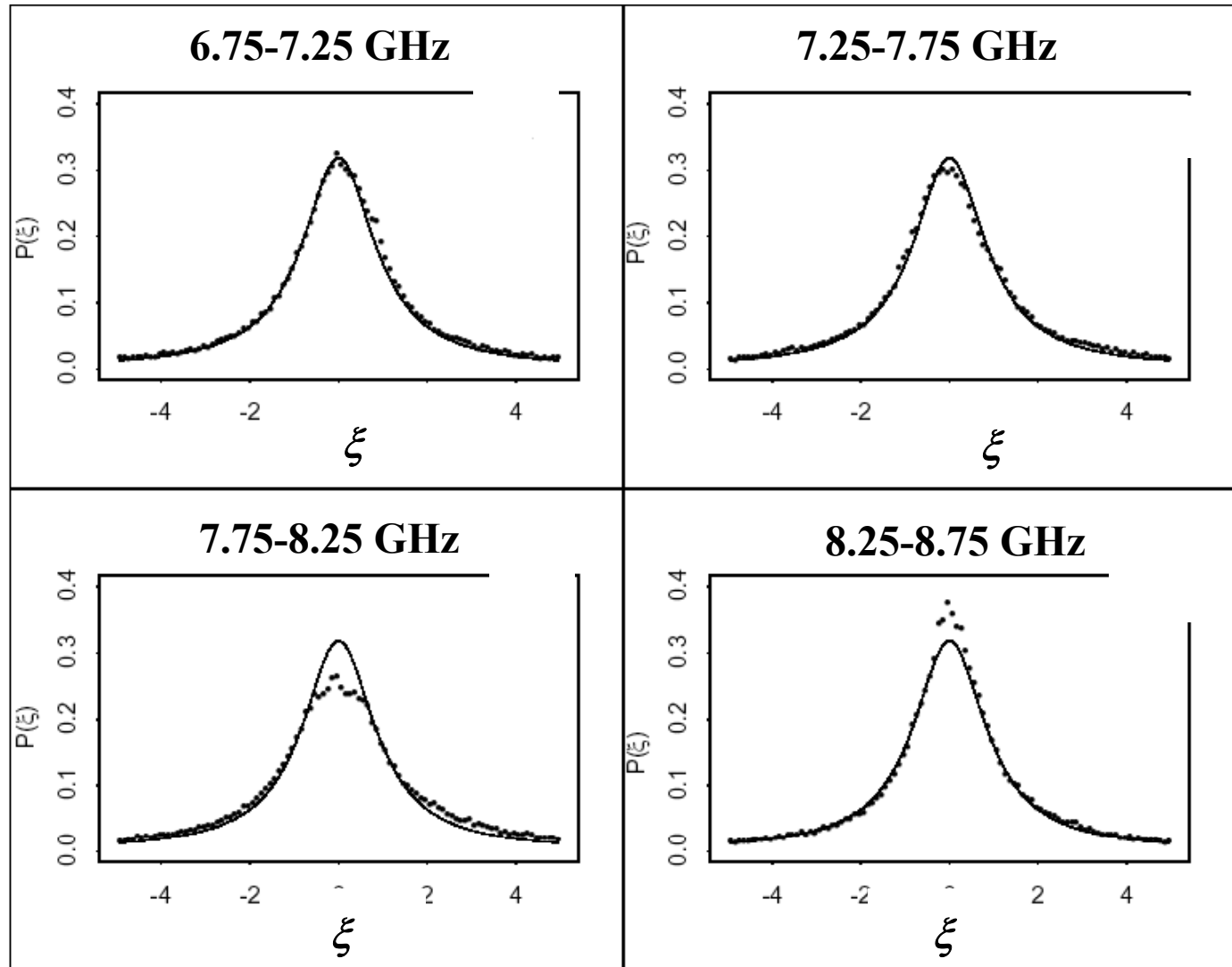




# Distribution of Fluctuating Cavity Impedance

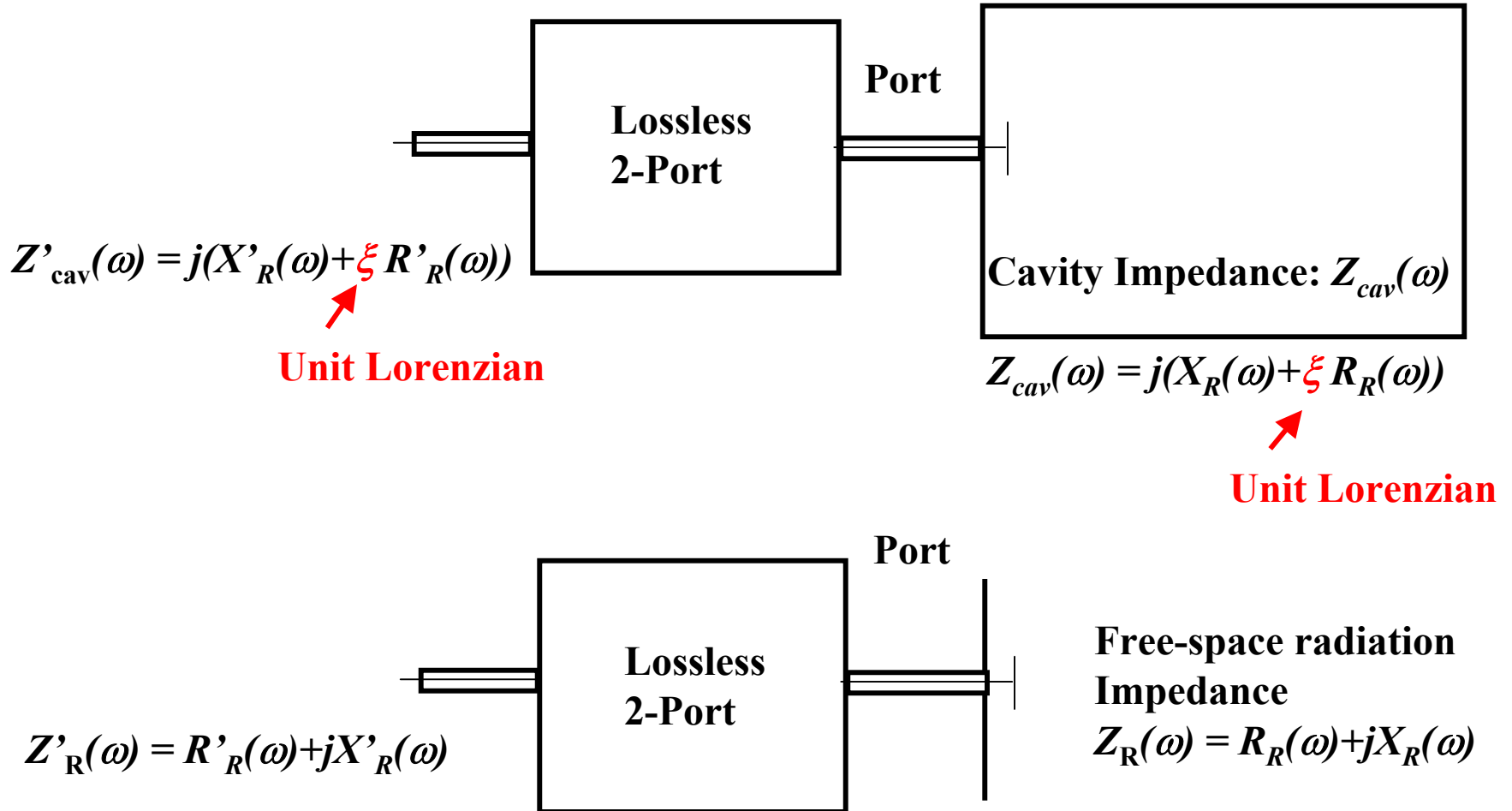
$$\xi = (X - X_R(\omega)) / R_R(\omega)$$

$R_{fs} \approx 35 \Omega$





# Impedance Transformation by Lossless Two-Port





# Properties of Lossless Two-Port Impedance

## Eigenvalues of $Z$ matrix

$$\det|Z - jXI| = 0$$

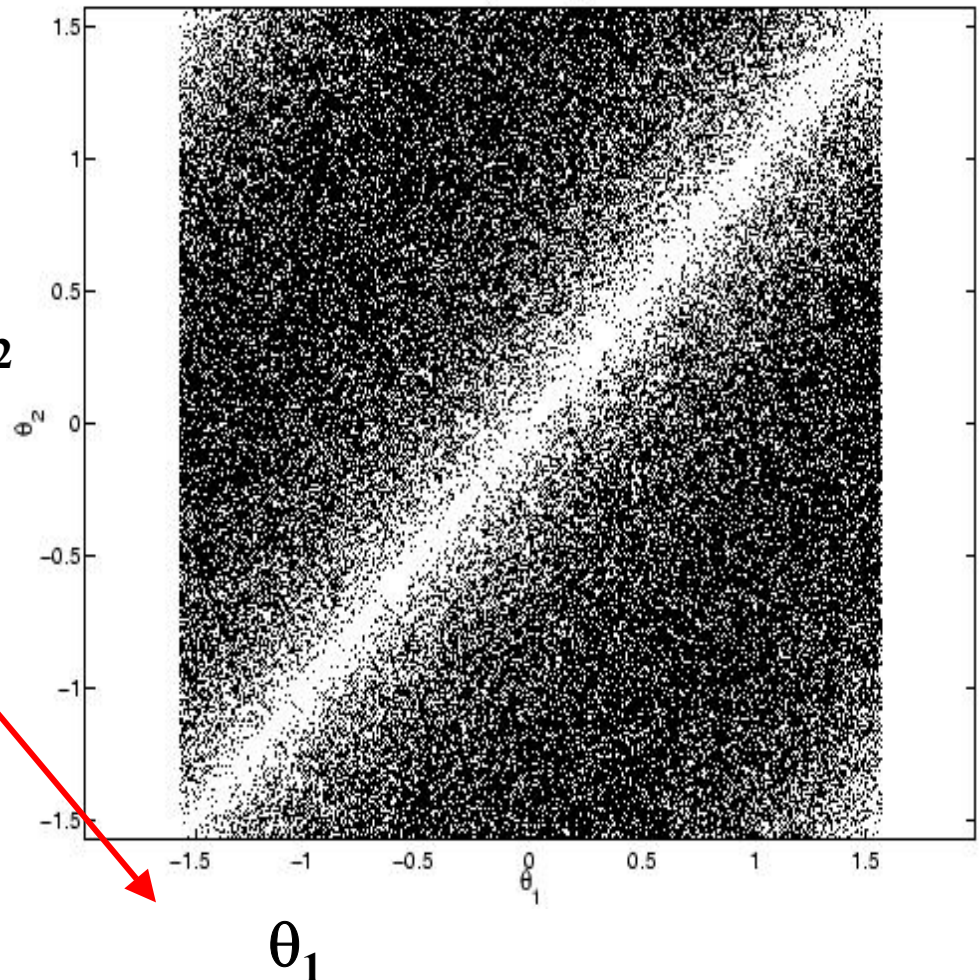
$$X_{1,2} = X_R + \xi_{1,2} R_R$$

$$\xi_{1,2} = \tan\left(\frac{\theta_{1,2}}{2}\right)$$

Individually  $\xi_{1,2}$  are  
Lorenzian distributed

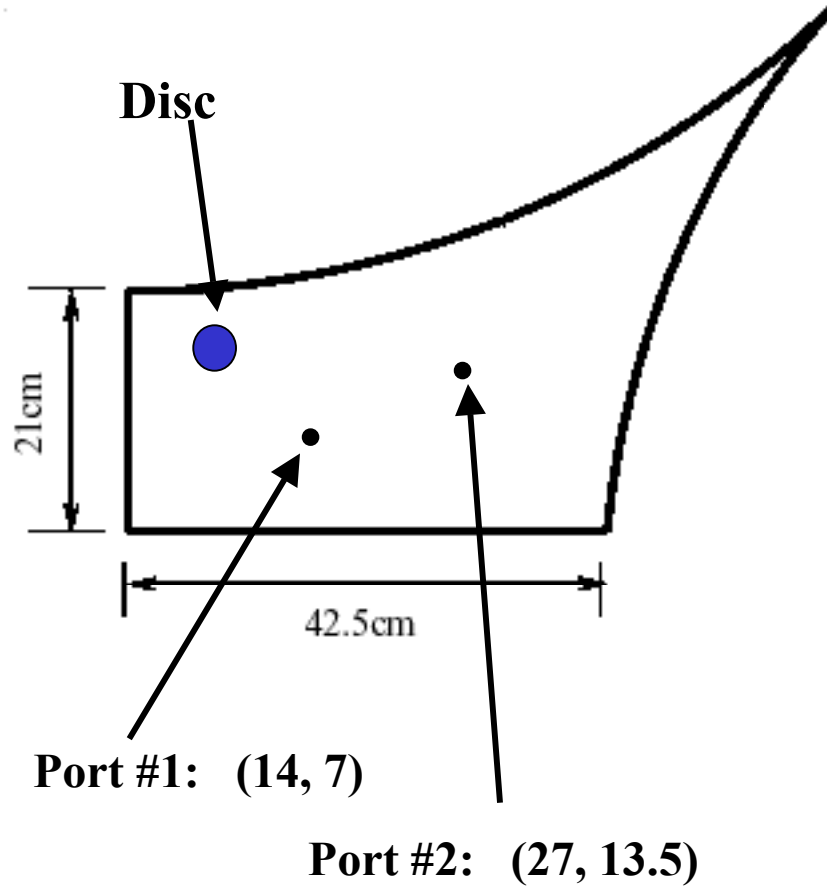
**Distributions same as  
In Random Matrix theory**

Joint Plot of  $\theta_1$  and  $\theta_2$  in TRS case

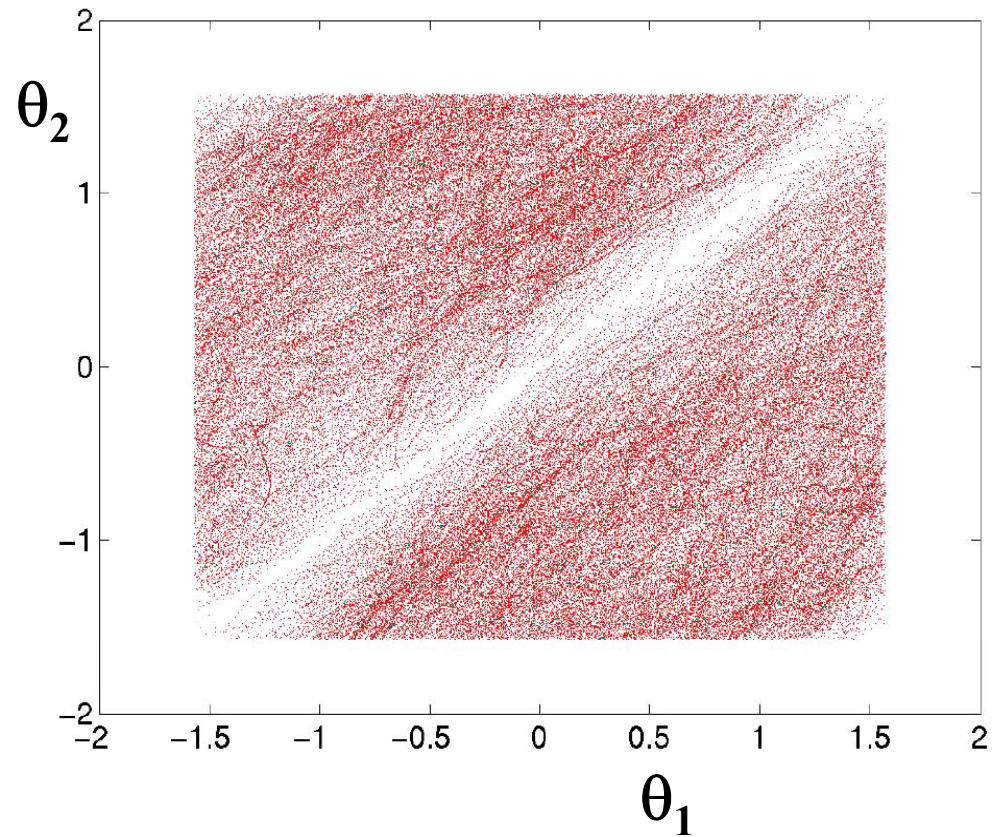




# HFSS Solution for Lossless 2-Port

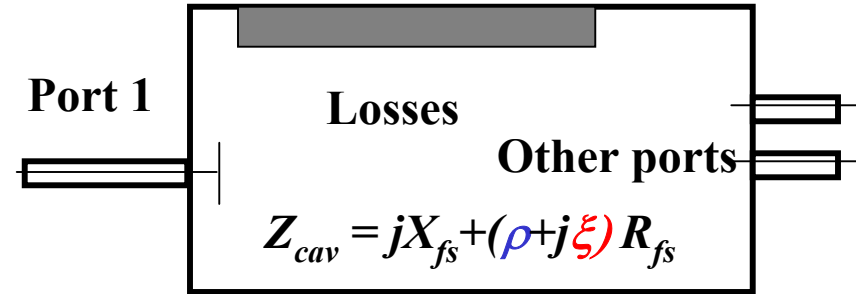


## Joint Pdf for $\theta_1$ and $\theta_2$



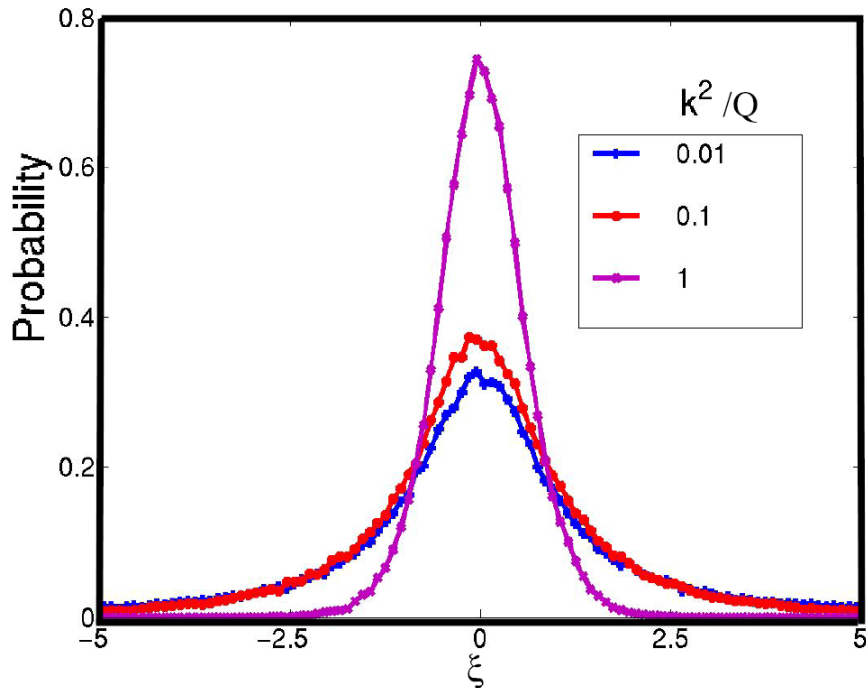


# Effect of Losses



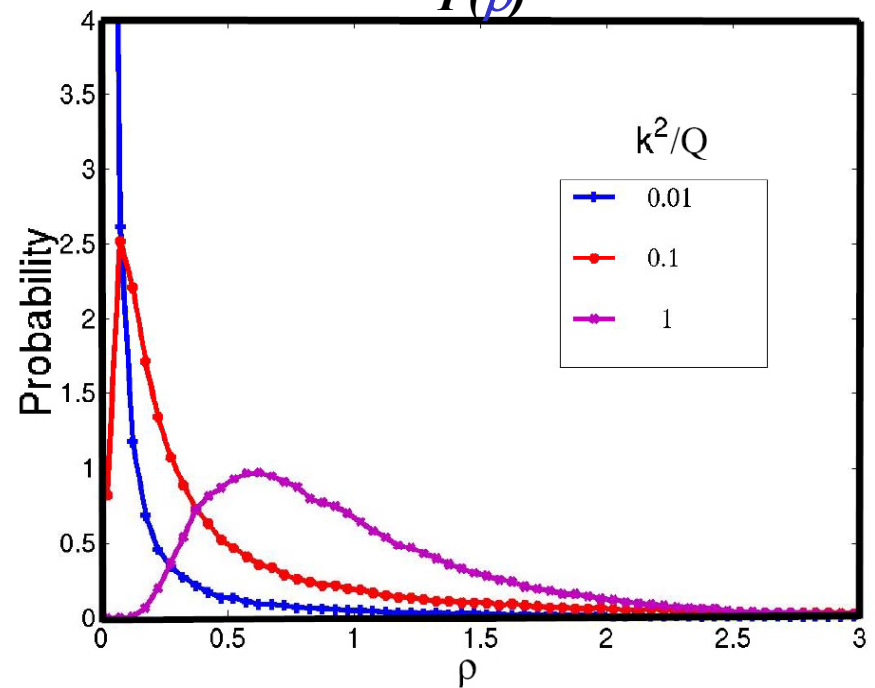
Distribution of reactance fluctuations

$$P(\xi)$$



Distribution of resistance fluctuations

$$P(\rho)$$



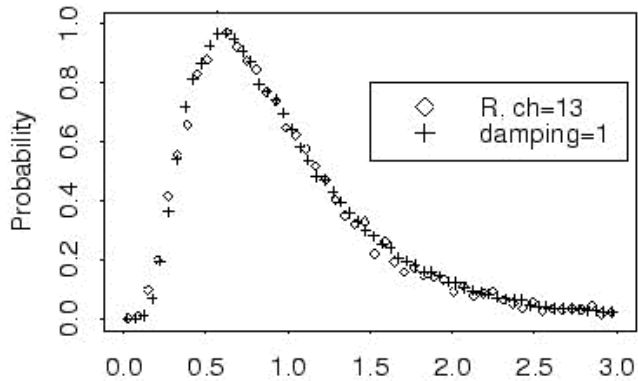


# Equivalence of Losses and Channels

$$Z_{cav} = jX_R + (\rho + j\xi)R_R$$

## Distribution of resistance fluctuations

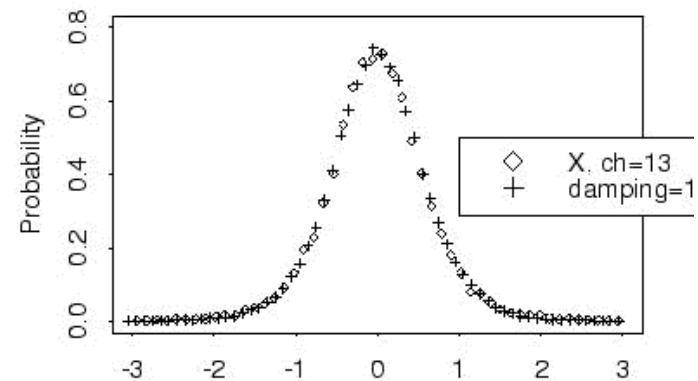
$P(\rho)$



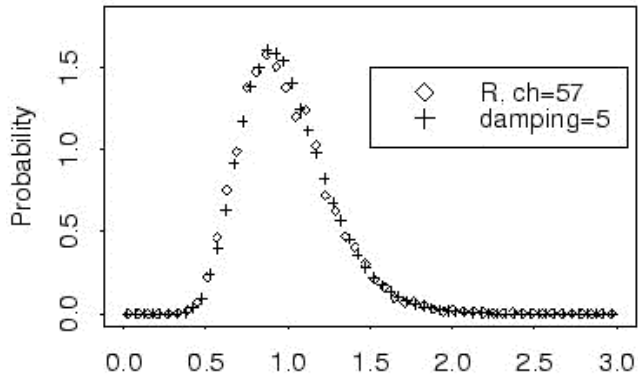
$\rho$

## Distribution of reactance fluctuations

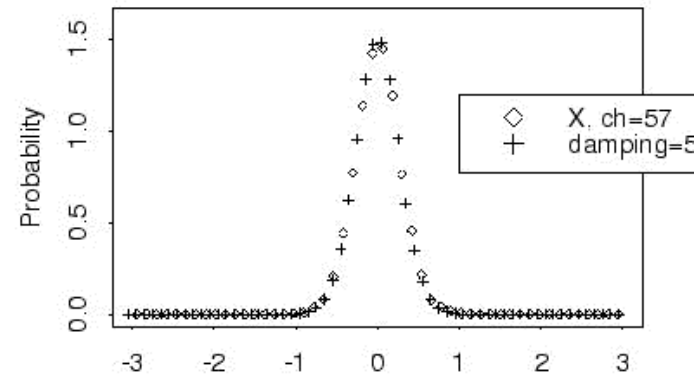
$P(\xi)$



$\xi$



$\rho$



$\xi$





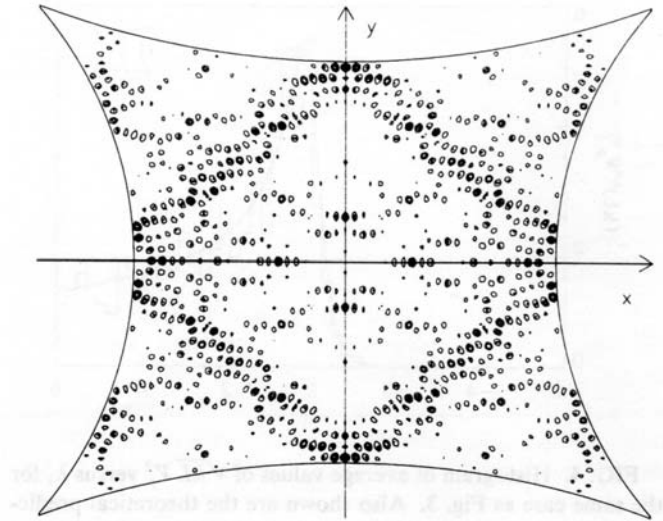
## Future Directions

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- Direct comparison of random coupling model with
  - random matrix theory ✓
  - HFSS solutions ✓
  - Experiment ✓
- Exploration of increasing number of coupling channels ✓
- Study losses in HFSS
- 3D examples
- Role of Scars on low period orbits
- Generalize to systems consisting of circuits and fields

# Role of Scars?

- Eigenfunctions that do not satisfy random plane wave assumption
- Scars are not treated by either random matrix or chaotic eigenfunction theory
- Semi-classical methods



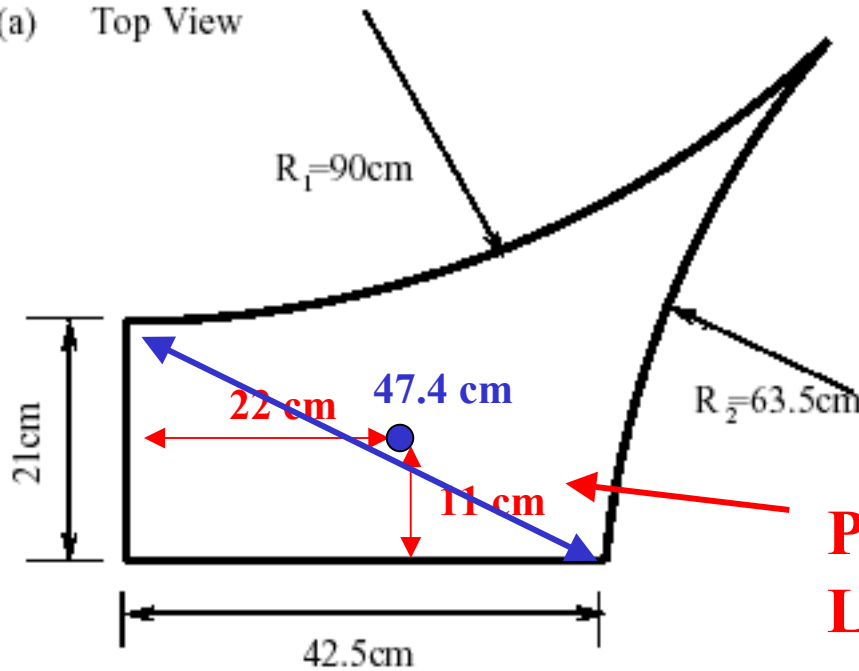
Bow-Tie with diamond scar





# Large Contribution from Periodic Ray Paths ?

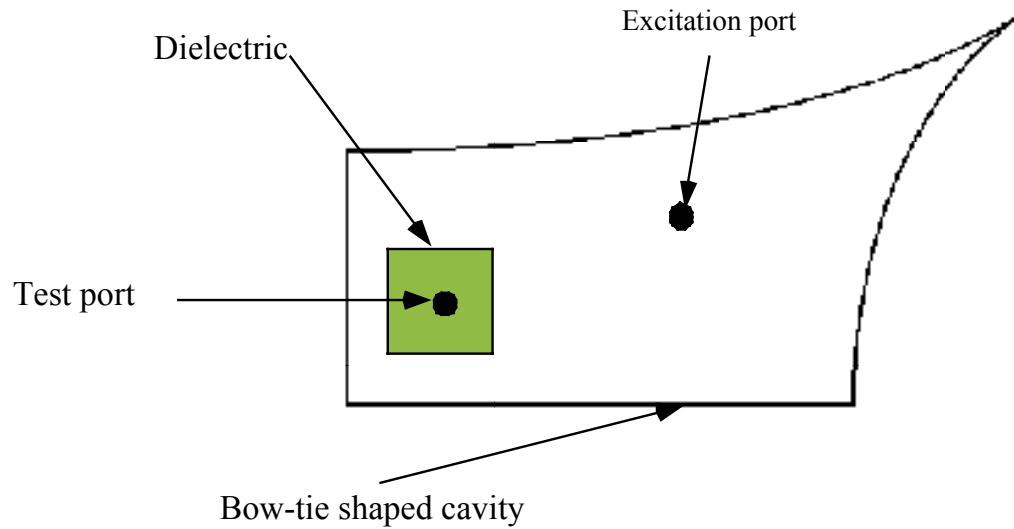
(a) Top View



**Possible strong reflections**  
 **$L = 94.8$  cm,  $\Delta f = .3$  GHz**



# Future Directions



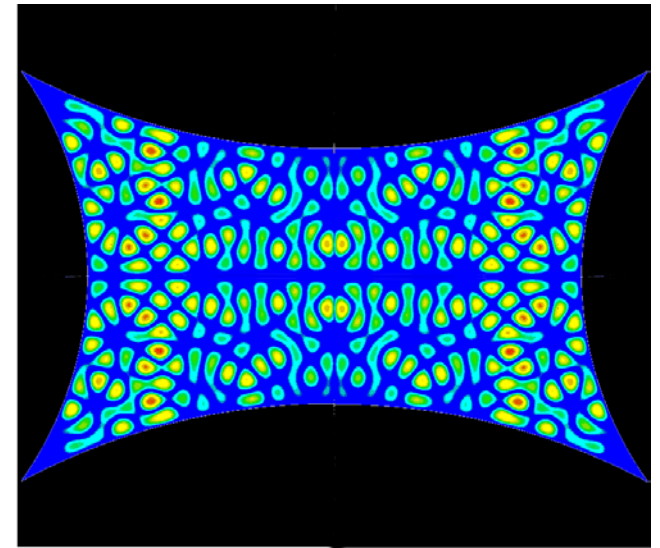
- Can be addressed
  - theoretically
  - numerically
  - experimentally

Features:

Ray splitting

Losses

Additional complications to be added later



HFSS simulation courtesy J. Rodgers